

# Bayesian inference of the missing resonances in $K^+ \Lambda$ photoproduction

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**Abstract.** A Regge-plus-resonance framework featuring consistent couplings for nucleon resonances up to spin  $J = 5/2$  is adopted to perform a Bayesian analysis of the world's  $\gamma p \rightarrow K^+ \Lambda$  data. It is concluded that the following nucleon resonances have the highest probability of decaying into the  $K^+ \Lambda$  channel:  $S_{11}(1535)$ ,  $S_{11}(1650)$ ,  $F_{15}(1680)$ ,  $P_{13}(1720)$ ,  $D_{13}(1900)$ ,  $P_{13}(1900)$ ,  $P_{11}(1900)$ , and  $F_{15}(2000)$ .

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## INTRODUCTION

The strange quark plays a peculiar role in Quantum Chromodynamics (QCD) as it is neither light nor heavy. A thorough knowledge of the nucleon-resonance ( $N^*$ ) content of open-strangeness production  $\gamma p \rightarrow K^+ \Lambda$  reactions would improve our understanding of the energy spectrum and the strong decay properties of baryons. The  $\gamma p \rightarrow K^+ \Lambda$  cross sections are of the order of  $\mu b$  and there are no outspoken structures in their measured energy dependence [1]. This points towards overlapping resonances and/or a dominant role for the background diagrams. Obviously, these observations complicate the extraction of the relevant  $N^*$  content from the data. In Ref. [2] a single-channel Regge-plus-resonance (RPR) model which cleanly separates the resonant and non-resonant contributions to  $\gamma p \rightarrow K^+ \Lambda$ , is proposed. The background is described in terms of Reggeized  $t$ -channel  $K^+(494)$  and  $K^{*+}(892)$  exchange and can be parametrized by three coupling strengths ( $g_{K^+ \Lambda p}$ ,  $G_{K^{*+} \Lambda p}^v$ ,  $G_{K^{*+} \Lambda p}^t$ ) and two phases (either 1 or  $\exp[-i\pi\alpha_K(t)]$ ) [2, 3, 4, 5]. The background parameters are constrained against the 262 data points from Jefferson lab with photon energies  $E_\gamma > 2.6$  GeV where no individual resonances are expected to contribute. The resonances are implemented with Feynman  $s$ -channel diagrams for  $N^*$  with  $J = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}$  and  $M_{N^*} \leq 2$  GeV. We adopt consistent couplings for the resonances [6]. These couplings are paramount for the analysis presented here, and require one parameter for each  $J = \frac{1}{2}$  resonance and two parameters for each  $N^*$  with  $J \geq \frac{3}{2}$ . In addition, there is a cutoff parameter in the hadronic form factor, which we assume identical for all  $N^*$ .

# BAYESIAN MODEL SELECTION AND THE RESONANT CONTENT OF $\gamma p \rightarrow K^+ \Lambda$

In order to determine the resonant content of the  $\gamma p \rightarrow K^+ \Lambda$  reaction in a statistically sound way, we have evaluated the  $N^*$  content of the RPR model against the world's data in a Bayesian analysis [3, 4]. The data comprise 3455 differential cross sections, 2241 single polarization observables (beam, target, and recoil) and, 452 double polarization observables (only beam-recoil data is published). In the analysis, 11 resonance candidates with masses up to 2 GeV are included. The spin-isospin quantum numbers and masses of the included  $N^*$  can be found in Fig. 1. We evaluate all possible combinations of the 11 candidate resonances. The best model is selected from the  $2^{11} = 2048$  model variants by calculating the Bayesian evidence  $\mathcal{Z}$  for each model against the world's  $p(\gamma, K^+) \Lambda$  data. A model's probability  $P(M|\{d_k\})$  is defined as the probability for a model  $M$  given the data  $\{d_k\}$ . Bayes' theorem allows one to connect this quantity to the evidence  $\mathcal{Z} = P(\{d_k\}|M)$ . The determination of  $\mathcal{Z}$  is computationally very expensive as it requires the likelihood for all possible values of the model parameters [3]. For the priors of the  $N^*$  parameters we adopt a uniform distribution whereby the upper bound is determined by naturalness arguments: no resonance is expected to generate strength which exceeds 5 times the measured total cross section. The absolute  $\mathcal{Z}_i$  values have little meaning, and only relative quantities  $\Delta \ln \mathcal{Z} \equiv \ln \mathcal{Z}_A / \mathcal{Z}_B$  can be interpreted with the aid of Jeffreys' scale [5]. We find that a model (coined RPR-2011) with 14  $N^*$  parameters provides the best evidence of describing the data. The RPR-2011 features the  $S_{11}(1535)$ ,  $S_{11}(1650)$ ,  $F_{15}(1680)$ ,  $P_{13}(1720)$ ,  $D_{13}(1900)$ ,  $P_{13}(1900)$ ,  $P_{11}(1900)$ , and  $F_{15}(2000)$  resonances. It is worth noticing that we find decisive evidence (measured according to Jeffreys' scale [3, 5]) that model variants with a larger amount of  $N^*$  parameters provide a worse description of the current data than RPR-2011. This proves that in a Bayesian analysis the most complex model does not necessarily emerge as the "best" model and that there is a cost for introducing additional model parameters.

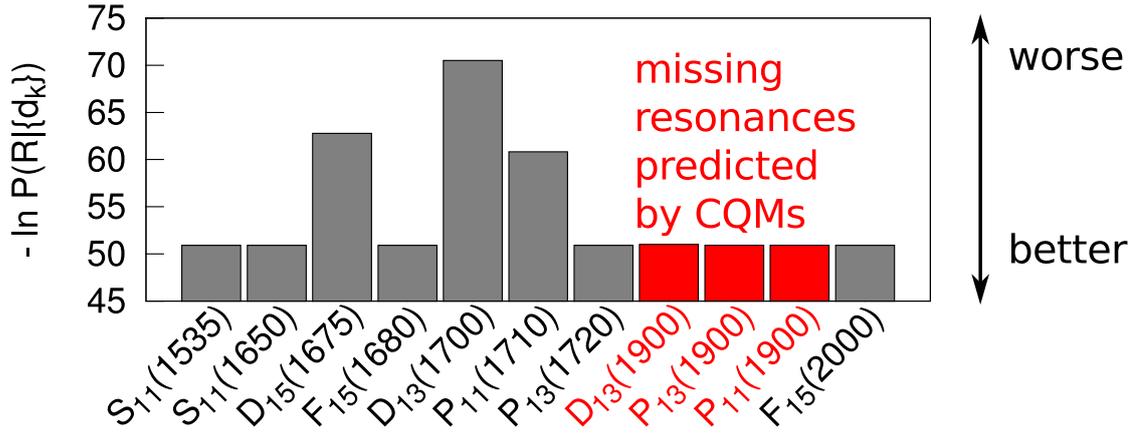
The probability of a resonance  $R$  given the  $\gamma p \rightarrow K^+ \Lambda$  data can be obtained by marginalizing over all possible models  $M_i$  which include a specific resonance  $R$

$$P(R|\{d_k\}) \Rightarrow \sum_{M_i|R \in M_i} P(M_i|\{d_k\}) = \sum_{M_i|R \in M_i} \underbrace{P(\{d_k\}|M_i)}_{\mathcal{Z}_i} \frac{P(M_i)}{P(\{d_k\})}. \quad (1)$$

Results are contained in Fig. 1 and confirm the resonance content of RPR-2011.

## CONCLUSION

The  $\gamma + p \rightarrow K^+ + \Lambda$  reaction is background dominated and has an overlapping  $N^*$  content which makes conventional isobar approaches and analysis schemes less appropriate. Bayesian methodology is a great tool for model selection under the condition that a moderate number of tunable parameters is involved. This confines the Bayesian methodology to a single-channel analysis. A Bayesian analysis within a coupled-channel framework is computationally prohibitive with the current computational resources. The RPR approach is a single-channel model for  $K^+ \Lambda$  photoproduction and provides an economical



**FIGURE 1.** The computed  $P(R|\{d_k\})$  as they have been defined in Eq. 1. There is decisive evidence for three resonances predicted by constituent quark models (CQM) and a mass of about 1900 MeV.

description from threshold up to  $E_\gamma^{\text{lab}} \leq 16$  GeV, the highest energy for which exclusive data are available. A Bayesian analysis of the world's data within the RPR framework points to 8  $N^*$ 's with a decay into  $K^+\Lambda$ . Thereby we find evidence for the  $P_{13}(1900)$  with an  $***$  overall status ( $***$  status for seen in  $K^+\Lambda$ ) in PDG [9], and for the  $P_{11}(1900)$ ,  $D_{13}(1900)$  which are “missing resonances”. We have evaluated the merits of the RPR model and have found that it has predictive power for  $ep \rightarrow e'K^+\Lambda$  [3], that it can predict the observables for kaon-hyperon production from the neutron [7], and that it provides a good elementary production operator for  $d(\gamma, K^0)Y$  processes [8].

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